## Exercises to Relativistic Quantum Field Theory — Sheet 7

— Prof. S. Dittmaier, Universität Freiburg, SS18 —

## Exercise 7.1 Wick theorem for bosonic fields (2 points)

The purpose of this exercise is to prove Wick's theorem for bosonic, real, free field operators  $\phi_i \equiv \phi_i(x_i)$ , which states that

$$T[\phi_1 \cdots \phi_n] = :\phi_1 \cdots \phi_n : + \sum_{\text{pairs } ij} :\phi_1 \cdots \phi_i \cdots \phi_j \cdots \phi_n : + \sum_{\text{double pairs } ij,kl} :\phi_1 \cdots \phi_i \cdots \phi_k \cdots \phi_j \cdots \phi_l \cdots \phi_n : + \dots$$

with the contractions representing propagators,  $\phi_i \phi_j = \langle 0|T[\phi_i \phi_j]|0\rangle$ . For n=2 the theorem has already been proven in Exercise 6.2. We organise the general proof in two steps. Without loss of generality we can assume that  $t_n = x_n^0$  is the smallest time variable, i.e.  $T[\phi_1 \cdots \phi_n] = T[\phi_1 \cdots \phi_{n-1}]\phi_n$ .

a) First prove the lemma

$$: \phi_1 \cdots \phi_{n-1} : \phi_n = : \phi_1 \cdots \phi_n : + \sum_{k=1}^{n-1} : \phi_1 \cdots \phi_k \cdots \phi_n : .$$

To this end, split  $\phi_n$  according to  $\phi_n = \phi_n^{(+)} + \phi_n^{(-)}$  into their positive and negative frequency parts  $\phi_n^{(\pm)}$ , i.e.  $\phi_n^{(+)}$  involves only annihilation operators and  $\phi_n^{(-)}$  only creation operators. For  $\phi_n^{(+)}$  the lemma is trivially verified, for  $\phi_n^{(-)}$  you can proceed via induction in n.

b) Argue that the lemma of a) trivially generalises to cases where contractions already appear inside the normal orderings, for example:

$$: \phi_1 \cdots \phi_i \cdots \phi_j \cdots \phi_{n-1} : \phi_n = : \phi_1 \cdots \phi_i \cdots \phi_j \cdots \phi_n : + \sum_{k=1}^{n-1} : \phi_1 \cdots \phi_i \cdots \phi_k \cdots \phi_j \cdots \phi_n : .$$

c) Prove Wick's theorem via induction in n using the result of b).

Please turn over!

Exercise 7.2 S-operator for two interacting scalar fields (1 point)

Consider a theory of a complex scalar field  $\phi$  (particle  $\phi$  and antiparticle  $\bar{\phi}$ ) and a real scalar field  $\Phi$  (particle  $\Phi$ ) with the Lagrangian density given by

$$\mathcal{L} = \frac{1}{2} (\partial_{\mu} \Phi)(\partial^{\mu} \Phi) - \frac{1}{2} M^2 \Phi^2 + (\partial_{\mu} \phi^{\dagger})(\partial^{\mu} \phi) - m^2 \phi^{\dagger} \phi + \mathcal{L}_{\text{int}},$$

where  $\mathcal{L}_{\text{int}} = \lambda \phi^{\dagger} \phi \Phi$ . Expand the S-operator,

$$S = T \exp\left(i \int d^4 x \, \mathcal{L}_{int}(x)\right),\,$$

up to order  $\lambda^2$  and use Wick's theorem to express the result in terms of propagators and normal-ordered products of fields. Note that the  $\lambda^n$  contribution can be written in the form

$$\frac{1}{n!} \int \mathrm{d}^4 x_1 \dots d^4 x_n : \dots :$$

Represent the result diagrammatically using the following notation:

• External lines:

$$\phi^{\dagger}(x) = \xrightarrow{x}, \quad \phi(x) = \xrightarrow{x}, \quad \Phi(x) = \cdots \xrightarrow{x}$$

• Internal lines:

$$\overrightarrow{\phi(x_1)}\overrightarrow{\phi^{\dagger}}(x_2) = \underbrace{x_1 \quad x_2}_{\bullet \bullet}, \qquad \overrightarrow{\Phi(x_1)}\overrightarrow{\Phi}(x_2) = \underbrace{x_1 \quad x_2}_{\bullet - \dots \bullet}$$

• Vertices:

$$i\lambda =$$