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Exercise 3.1 Virial theorem (1 point)

Consider a spinless particle of mass m in a potential $V(\vec{x})$. Show the relation

$$2\langle T \rangle_{\phi_n} = \langle \hat{\vec{x}} \cdot \nabla V \rangle_{\phi_n} \tag{1}$$

for expectation values $\langle ... \rangle_{\phi_n}$ in (stationary) energy eigenstates $|\phi_n\rangle$, where $T = \hat{\vec{p}}^2/(2m)$ is the operator for the kinetic energy of the particle. What does this relation imply for the expectation values $\langle T \rangle_{\phi_n}$ and $\langle V \rangle_{\phi_n}$ for central potentials of the type $V(\vec{x}) \propto r^s$, where $r = |\vec{x}|$?

(Hint: Evaluate $\langle [\hat{\vec{x}} \cdot \hat{\vec{p}}, \hat{H}] \rangle_{\phi_n}$ is two different ways.)

Exercise 3.2 Electromagnetic gauge transformations (3 points)

The Hamiltonian of a non-relativistic, spinless particle of mass m and electric charge q in a classical electromagnetic field is given by

$$\hat{H}(\hat{\vec{x}}, \hat{\vec{p}}) = \frac{1}{2m} \left(\hat{\vec{p}} - q \vec{A}(\hat{\vec{x}}, t) \right)^2 + q \Phi(\hat{\vec{x}}, t), \tag{2}$$

where $\vec{A}(\vec{x},t)$ and $\Phi(\vec{x},t)$ are the classical vector and scalar potentials of the electromagnetic field, respectively. Here, $\hat{\vec{x}}$ is the usual position operator and $\hat{\vec{p}}$ its canonical conjugate momentum.

a) The electric and magnetic field strengths $\vec{E} = -\nabla \Phi - \vec{A}$ and $\vec{B} = \nabla \times \vec{A}$ are invariant under the gauge transformation:

$$\vec{A} \rightarrow \vec{A}' = \vec{A} + \nabla \chi, \qquad \Phi \rightarrow \Phi' = \Phi - \dot{\chi},$$
 (3)

where $\chi = \chi(\vec{x}, t)$ is an arbitrary real function of space and time. Show that the Hamilton operator transforms as

$$\hat{H} \rightarrow \hat{H}' = U\hat{H}U^{\dagger} + i\hbar\dot{U}U^{\dagger},$$
 (4)

with the operator $U(\hat{\vec{x}},t) = \exp\{iq\chi(\hat{\vec{x}},t)/\hbar\}$. Is U unitary?

- b) Show that $|\psi'(t)\rangle = U|\psi(t)\rangle$ obeys the time-dependent Schrödinger equation with Hamiltonian \hat{H}' if $|\psi(t)\rangle$ obeys the Schrödinger equation with Hamiltonian \hat{H} .
- c) Identify the operator $m\hat{\vec{v}}$ corresponding to the cartesian momentum $m\dot{\vec{x}}$ as the operator that produces the expectation value $m\frac{d}{dt}\langle\hat{\vec{x}}\rangle$. What are the commutators $[\hat{x}_k, m\hat{v}_l]$? Consider the two momentum expectation values $\langle\hat{\vec{p}}\rangle$ and $\langle m\hat{\vec{v}}\rangle$. Which of the two is invariant under gauge transformations (3)?

Exercise 3.3 Discrete rotations in two dimensions – the cyclic groups (3 points)

For a given natural number n, consider the group of discrete rotations about integer multiples of the angle $\frac{2\pi}{n}$ in two dimensions, which are represented by the matrices

$$R(\phi_k) = \begin{pmatrix} \cos \phi_k & -\sin \phi_k \\ \sin \phi_k & \cos \phi_k \end{pmatrix}, \qquad \phi_k = \frac{2\pi k}{n} \qquad k = 0, 1, \dots, n - 1.$$
 (5)

These matrices define a two-dimensional representation of the cyclic group C_n .

- a) Determine the similarity transformation that reduces the representation (5) to two irreducible representations.
- b) The "regular representation" of C_3 (and analogously for C_n) is defined by the following three matrices:

$$D(e) = \mathbf{1}, \qquad D(g) = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}, \qquad D(g^2) = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}, \tag{6}$$

where g is the *generating element* of the group. Similarly to a) fully reduce this representation.

(Hint: Introduce $\epsilon = e^{2\pi i/3}$, with $\epsilon^2 = \epsilon^*$, $1 + \epsilon + \epsilon^2 = 0$.)

c) C_n possesses n one-dimensional inequivalent representations. Guess them from the pattern observed for C_3 in b).

Exercise 3.4 Parity operator (2 points)

The parity operator \mathcal{P} is linear and defined by the following action on position eigenstates:

$$\mathcal{P}|\vec{x}\rangle = |-\vec{x}\rangle. \tag{7}$$

(Spin will not be considered in this exercise.)

- a) Show that \mathcal{P} is unitary and acts on position-space wave functions as $\mathcal{P}\psi(\vec{x}) = \psi(-\vec{x})$. Derive the operators $\hat{\vec{x}}'$, $\hat{\vec{p}}'$, $\hat{\vec{L}}'$, where $A' = \mathcal{P} A \mathcal{P}^{-1}$ is the parity-transformed version of an operator A. Here $\hat{\vec{L}} = \hat{\vec{x}} \times \hat{\vec{p}}$ is the usual orbital angular momentum of a single particle.
- b) Derive the parity-transformed operators of the electromagnetic potentials and field strengths $\vec{A}'(\hat{\vec{x}},t), \; \Phi'(\hat{\vec{x}},t), \; \vec{E}'(\hat{\vec{x}},t), \; \vec{B}'(\hat{\vec{x}},t), \; \text{upon analyzing Maxwell's equations}$ and the fact that electric charges do not change under parity transformations.