Exercises on Supersymmetry S	Sheet 9	
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**Exercise 19** Auxiliary relations for conventional gauge transformations (3 points)

A non-abelian local gauge transformation can represented using the unitary matrix

$$U\left(\omega^{A}(x)\right) = \exp\left\{-igT^{A}\omega^{A}(x)\right\},\tag{1}$$

where  $T^A$  denotes the generators of the gauge group and  $\omega^A$  are the group parameters that are real functions of space-time  $x^{\mu}$ . Show the relation

$$\exp\left\{-igT^{A}\omega^{A}(\bar{y})\right\} = U + i(\theta\bar{\sigma}^{\mu}\bar{\theta})(\partial_{\mu}U) - \frac{1}{4}(\theta\theta)(\bar{\theta}\bar{\theta})(\Box U),\tag{2}$$

which is an expansion of the generalized exponential as a function of coordinates  $\bar{y}^{\mu} = x^{\mu} + i\theta\bar{\sigma}^{\mu}\bar{\theta}$ . The exponentials on the r.h.s. of Eq. (2) are functions of x, as defined in Eq. (1).

## **Exercise 20** Superfield of the non-abelian field-strength tensor (8 points)

In Wess–Zumino gauge the components of a vector superfield of a non-abelian gauge symmetry reads

$$V^{A}(x,\theta,\bar{\theta})\Big|_{WZ} = -(\theta\bar{\sigma}^{\mu}\bar{\theta})A^{A}_{\mu}(x) - i(\theta\theta)\bar{\theta}\bar{\lambda}^{A}(x) + i(\bar{\theta}\bar{\theta})\theta\lambda^{A}(x) - \frac{1}{2}(\theta\theta)(\bar{\theta}\bar{\theta})D^{A}(x), \quad (3)$$

with  $A = 1, \ldots, \dim(G)$ . The corresponding (SUSY gauge-covariant) superfield of the non-abelian field-strength tensor is defined as

$$TW_a(x,\theta,\bar{\theta}) \equiv T^A W_a^A(x,\theta,\bar{\theta})$$
  
=  $\frac{1}{8g} (\bar{\mathcal{D}}\bar{\mathcal{D}}) \exp\{2gTV(x,\theta,\bar{\theta})\} \mathcal{D}_a \exp\{-2gTV(x,\theta,\bar{\theta})\},$  (4)

with  $T^A$  being the generators of the gauge group G which satisfy  $[T^A, T^B] = iC^{ABC}T^C$ . Calculate the component form of  $TW_a$  in terms of the coordinates  $y^{\mu} = x^{\mu} - i\theta \bar{\sigma}^{\mu} \bar{\theta}, \eta = \theta$ , and  $\bar{\eta} = \bar{\theta}$ :

$$TW_a(y,\eta) = T\lambda_a(y) + i\eta_a(TD(y)) - i(\eta\eta)\bar{\sigma}^{\mu}_{ab}\left(D_{\mathrm{adj},\mu}\bar{\lambda}^b(y)\right) + \frac{1}{2}\bar{\sigma}^{\mu}_{ab}\sigma^{bc,\nu}\eta_c\left(TF_{\mu\nu}(y)\right).$$
(5)

The matrix  $D_{\mathrm{adj},\mu}$  denotes covariant derivative in adjoint representation (of G),

$$(D_{\mathrm{adj},\mu})^{BC} = \delta^{BC} \partial_{\mu} + g C^{ABC} A^{A}_{\mu}, \tag{6}$$

and  $F^A_{\mu\nu}$  are the components of the non-abelian field-strength tensor,

$$F^A_{\mu\nu} = \partial_\mu A^A_\nu - \partial_\nu A^A_\mu - g C^{ABC} A^B_\mu A^C_\nu.$$
<sup>(7)</sup>