

Exercise 4.1 (4 points) *Massive gauge-boson propagator in R_ξ gauge*

The propagator $D_\xi^{\mu\nu}(x)$ of a massive gauge boson with mass M is defined by

$$\left[g_{\mu\nu} (\partial^2 + M^2) + \left(\frac{1}{\xi} - 1 \right) \partial_\mu \partial_\nu \right] D_\xi^{\nu\rho}(x) = \delta_\mu^\rho \delta(x) .$$

- a) Calculate the Fourier-transformed $\tilde{D}_\xi^{\mu\nu}(q)$ of the propagator by inserting

$$D_\xi^{\mu\nu}(x) = \int \frac{d^4q}{(2\pi)^4} \exp\{iqx\} \tilde{D}_\xi^{\mu\nu}(q)$$

into the differential equation given above. Make use of the decomposition of $\tilde{D}_\xi^{\mu\nu}(q)$ into transverse and longitudinal parts, $\tilde{D}_{T,\xi}(q)$ and $\tilde{D}_{L,\xi}(q)$, respectively, with

$$\tilde{D}_\xi^{\mu\nu}(q) = \tilde{D}_{T,\xi}(q) \left(g^{\mu\nu} - \frac{q^\mu q^\nu}{q^2} \right) + \tilde{D}_{L,\xi}(q) \frac{q^\mu q^\nu}{q^2} .$$

Determine $\tilde{D}_\xi^{\mu\nu}(q)$ in the limits $\xi \rightarrow 0$, $\xi \rightarrow 1$, and $\xi \rightarrow \infty$.

- b) Given the generating functional

$$Z_0[J_\mu] = \frac{1}{N} \int \mathcal{D}A^\mu \exp \left\{ i \int d^4x [\mathcal{L}_0 + J_\mu A^\mu] \right\}$$

with

$$Z_0[0] = 1 , \quad \mathcal{L}_0 = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} M^2 A_\mu A^\mu - \frac{1}{2\xi} (\partial A)^2 ,$$

show that

$$Z_0[J_\mu] = \exp \left\{ +\frac{1}{2} \int d^4x \int d^4x' iJ_\mu(x) iD_\xi^{\mu\nu}(x-x') iJ_\nu(x') \right\} .$$

Please turn over!

Exercise 4.2 (3 points) *Two-point Green's function of the ϕ^4 theory*

The interaction part of the Lagrangian of a ϕ^4 theory with a single, real scalar field ϕ is given as

$$\mathcal{L}_I = -\frac{g}{4!}\phi^4.$$

Starting from the generating functionals, calculate the Green's function $G^{\phi\phi}(x_1, x_2)$ and the connected Green's function $G_{\text{con}}^{\phi\phi}(x_1, x_2)$ up to order $\mathcal{O}(g)$ and draw diagrams representing the resulting terms.

Exercise 4.3 (3 points) *Equation of motion for Green's functions*

Consider a quantum field theory of a real scalar field $\phi(x)$ with the Lagrangian $\mathcal{L}(\phi) = \mathcal{L}_0(\phi) + \mathcal{L}_I(\phi)$ where the free part is given by $\mathcal{L}_0(\phi) = -\frac{1}{2}\phi(\partial^2 + m^2)\phi$ and the interaction part $\mathcal{L}_I(\phi)$ is not further specified.

- a) Verify explicitly that the free generating functional

$$Z_0[J] = \exp\left\{+\frac{1}{2}\int d^4x \int d^4x' iJ(x)i\Delta_F(x-x')iJ(x')\right\}$$

fulfills the following equation of motion

$$\left[\frac{\delta\mathcal{L}_0}{\delta\phi}\left(\frac{\delta}{i\delta J(x)}\right) + J(x)\right]Z_0[J] = 0.$$

- b) Starting with this equation, derive the equations of motion for the free two- and four-point functions, $G_0^{\phi\phi}(x_1, x_2)$ and $G_0^{\phi\phi\phi\phi}(x_1, x_2, x_3, x_4)$ by taking the functional derivative.
- c) By explicitly inserting the generating functional

$$Z[J] = \exp\left\{i\int d^4y \mathcal{L}_I\left(\frac{\delta}{i\delta J(x)}\right)\right\}Z_0[J],$$

show that in a theory with interactions the following equation of motions hold,

$$\left[\frac{\delta\mathcal{L}}{\delta\phi}\left(\frac{\delta}{i\delta J(x)}\right) + J(x)\right]Z[J] = 0.$$

Use (and prove) the commutator relation

$$\left[\exp\left\{i\int d^4y \mathcal{L}_I\left(\frac{\delta}{i\delta J(x)}\right)\right\}, J(x)\right] = \frac{\delta\mathcal{L}_I}{\delta\phi}\left(\frac{\delta}{i\delta J(x)}\right)\exp\left\{i\int d^4y \mathcal{L}_I\left(\frac{\delta}{i\delta J(x)}\right)\right\}.$$