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**Exercises to Relativistic Quantum Field Theory — Sheet 8**

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**Exercise 8.1** *S-operator for two interacting scalar fields (cont'd)* (1 point)

Consider again the field theory of a complex scalar field  $\phi$  (particle  $\phi$  and antiparticle  $\bar{\phi}$ ) and a real scalar field  $\Phi$  (particle  $\Phi$ ) from Exercise 7.2 with the Lagrangian density

$$\mathcal{L} = \frac{1}{2}(\partial_\mu\Phi)(\partial^\mu\Phi) - \frac{1}{2}M^2\Phi^2 + (\partial_\mu\phi^\dagger)(\partial^\mu\phi) - m^2\phi^\dagger\phi + \mathcal{L}_{\text{int}},$$

where  $\mathcal{L}_{\text{int}} = \lambda\phi^\dagger\phi\Phi$ , and make use of the perturbative expansion worked out there.

- Calculate the  $S$ -matrix element  $S_{fi} = \langle f|S|i\rangle$  in lowest order between the initial state  $|i\rangle = a_\Phi^\dagger(k)|0\rangle$  and the final state  $|f\rangle = a_\phi^\dagger(p_1)b_\phi^\dagger(p_2)|0\rangle$ , where  $a_\Phi^\dagger(q)$ ,  $a_\phi^\dagger(q)$ ,  $b_\phi^\dagger(q)$  are the creation operators of the particles  $\Phi$ ,  $\phi$ , and  $\bar{\phi}$ .
- Assuming  $M > 2m$ , calculate the lowest-order decay width

$$\Gamma_{\Phi \rightarrow \phi\bar{\phi}} = \frac{1}{2M} \int d\Phi_2 |\mathcal{M}_{fi}|^2$$

for the decay  $\Phi \rightarrow \phi\bar{\phi}$ , where  $\Phi_2$  is the 2-particle phase space of the final state (see Exercise 5.2) and the transition matrix element  $\mathcal{M}_{fi}$  is related to  $S_{fi}$  by

$$S_{fi} = (2\pi)^4 \delta(k - p_1 - p_2) i\mathcal{M}_{fi}.$$

**Exercise 8.2** *Fundamental representations of the Lorentz group* (1 point)

The general form of Lorentz transformations in the two fundamental representations of the Lorentz group is given by

$$\Lambda_R = \exp\left(-\frac{i}{2} \sum_{k=1}^3 (\phi_k + i\nu_k) \sigma^k\right), \quad \Lambda_L = \exp\left(-\frac{i}{2} \sum_{k=1}^3 (\phi_k - i\nu_k) \sigma^k\right)$$

with the real group parameters  $\phi_k$ ,  $\nu_k$  and the Pauli matrices  $\sigma^k$ .

- Show that  $\Lambda_R^\dagger = \Lambda_L^{-1}$  and  $\Lambda_L^\dagger = \Lambda_R^{-1}$ .
- Show that  $\det(\Lambda_R) = \det(\Lambda_L) = 1$  using  $\det(\exp\{A\}) = \exp\{\text{Tr}(A)\}$  for a matrix  $A$ .
- Which transformations are characterized by  $\Lambda_{R/L}^\dagger = \Lambda_{R/L}$ , which by  $\Lambda_{R/L}^\dagger = \Lambda_{R/L}^{-1}$ ?

- Derive  $\Lambda_R$  and  $\Lambda_L$  for a pure boost in the direction  $\vec{e} = \begin{pmatrix} \cos\varphi \sin\theta \\ \sin\varphi \sin\theta \\ \cos\theta \end{pmatrix}$  with  $\vec{\nu} = \nu\vec{e}$ ,  $\vec{\phi} = 0$  and for a pure rotation around the axis  $\vec{e}$  with  $\vec{\phi} = \phi\vec{e}$ ,  $\vec{\nu} = 0$ .

*Please turn over!*

**Exercise 8.3** *Connection between  $\Lambda_R$ ,  $\Lambda_L$ , and  $\Lambda^\mu{}_\nu$*  (1 point)

The general matrix representing a Lorentz transformation of a four-vector is given by

$$\Lambda^\mu{}_\nu = \exp\left(-\frac{i}{2}\omega_{\alpha\beta}M^{\alpha\beta}\right)^\mu{}_\nu, \quad (M^{\alpha\beta})^\mu{}_\nu = i(g^{\alpha\mu}g^\beta{}_\nu - g^{\beta\mu}g^\alpha{}_\nu)$$

with the antisymmetric parameters  $\omega_{jk} = \epsilon_{jkl}\phi_l$  and  $\omega_{0j} = -\omega_{j0} = \nu_j$ . The connection between  $\Lambda_R$ ,  $\Lambda_L$  (see Exercise 8.1) and  $\Lambda^\mu{}_\nu$  is

$$\Lambda_R^\dagger \sigma^\mu \Lambda_R = \Lambda^\mu{}_\nu \sigma^\nu, \quad \Lambda_L^\dagger \bar{\sigma}^\mu \Lambda_L = \Lambda^\mu{}_\nu \bar{\sigma}^\nu,$$

where  $\sigma^\mu = (\mathbb{1}, \sigma^1, \sigma^2, \sigma^3)$  and  $\bar{\sigma}^\mu = (\mathbb{1}, -\sigma^1, -\sigma^2, -\sigma^3)$ . Verify these relations for infinitesimal transformations with the parameters  $\delta\phi_k, \delta\nu_k$ .