

**Exercise 22**      *Sfermion sector of the MSSM*      (5 points)

Assuming “minimal flavour violation” one obtains the following mass terms for sfermions  $\tilde{f}$ ,

$$\mathcal{L}_{\text{mass},\tilde{f}} = - \begin{pmatrix} \hat{f}_L^\dagger & \hat{f}_R^\dagger \end{pmatrix} Z_{\tilde{f}} \begin{pmatrix} \hat{f}_L \\ \hat{f}_R \end{pmatrix}, \quad Z_{\tilde{f}} = \begin{pmatrix} m_f^2 + M_{\tilde{f}}^{\text{LL}} & m_f (M_{\tilde{f}}^{\text{LR}})^* \\ m_f M_{\tilde{f}}^{\text{LR}} & m_f^2 + M_{\tilde{f}}^{\text{RR}} \end{pmatrix}, \quad (1)$$

where  $\hat{f}_{L/R}$  are the sfermion fields that correspond to fermion fields  $\Psi_f$  in the basis of fermion mass eigenstates. The coefficients of the mass matrix  $Z_{\tilde{f}}$  are given as

$$M_{\tilde{f}}^{\text{LL}} = M_Z^2 \cos(2\beta) (I_{w,f_L}^3 - Q_f s_w^2) + M_{\tilde{F}_L}^2, \quad I_{w,f_L}^3 = \pm \frac{1}{2}, \quad (2a)$$

$$M_{\tilde{f}}^{\text{LR}} = A_f - \mu^* \begin{cases} \cot \beta, & I_{w,f_L}^3 = +\frac{1}{2} \\ \tan \beta, & I_{w,f_L}^3 = -\frac{1}{2} \end{cases}, \quad (2b)$$

$$M_{\tilde{f}}^{\text{RR}} = M_Z^2 \cos(2\beta) Q_f s_w^2 + M_{\tilde{F}_R}^2. \quad (2c)$$

Diagonalization of the mass matrix  $Z_{\tilde{f}}$  yields

$$\mathcal{L}_{\text{mass},\tilde{f}} = - \sum_{k=1,2} m_{\tilde{f}_k}^2 \tilde{f}_k^\dagger \tilde{f}_k, \quad (3)$$

with new sfermion fields

$$\begin{pmatrix} \tilde{f}_1 \\ \tilde{f}_2 \end{pmatrix} = U_{\tilde{f}} \begin{pmatrix} \hat{f}_L \\ \hat{f}_R \end{pmatrix}, \quad U_{\tilde{f}} = \begin{pmatrix} \cos \theta_{\tilde{f}} & \sin \theta_{\tilde{f}} \\ -\sin \theta_{\tilde{f}} & \cos \theta_{\tilde{f}} \end{pmatrix}, \quad (4)$$

where  $\theta_{\tilde{f}}$  is a real mixing angle. Calculate the mass values  $m_{\tilde{f}_k}^2$  (by definition  $m_{\tilde{f}_1} \leq m_{\tilde{f}_2}$ ) and show the following useful relations,

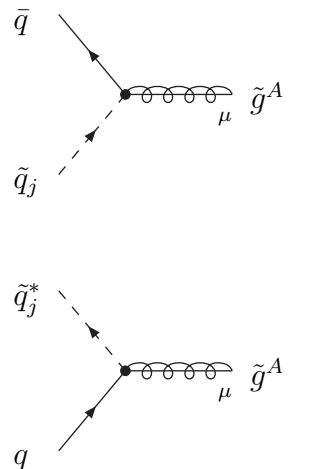
$$\cos(2\theta_{\tilde{f}}) = \frac{M_{\tilde{f}}^{\text{LL}} - M_{\tilde{f}}^{\text{RR}}}{m_{\tilde{f}_1}^2 - m_{\tilde{f}_2}^2}, \quad \sin(2\theta_{\tilde{f}}) = \frac{2m_f [A_f - \mu \{\cot \beta, \tan \beta\}]}{m_{\tilde{f}_1}^2 - m_{\tilde{f}_2}^2}, \quad (5)$$

with  $\{\cot \beta, \tan \beta\}$  denoting the alternatives for  $I_{w,f_L}^3 = \pm \frac{1}{2}$ . What happens in the limiting cases of

- a) small fermion mass  $m_f$ ,
- b) large SUSY parameters  $M_{\tilde{F}_L}$ ,  $M_{\tilde{F}_R}$ ,  $A_f$ , and  $\mu$ , which should scale proportional to a large scale  $M_{\text{SUSY}}$ ?

**Exercise 23** *SUSY QCD correction to the quark self-energy* (5 bonus points)

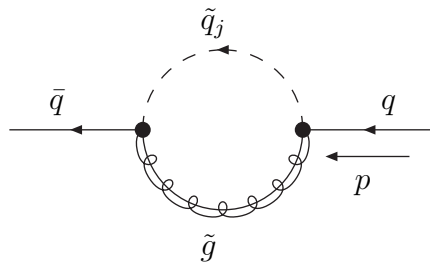
The Feynman rules for the quark–squark–gluon interaction are as follows,



$$i\sqrt{2}g_s \frac{\lambda^A}{2} (-U_{\tilde{q},j1}^* \omega_+ + U_{\tilde{q},j2}^* \omega_-)$$

$$i\sqrt{2}g_s \frac{\lambda^A}{2} (U_{\tilde{q},j2} \omega_+ - U_{\tilde{q},j1} \omega_-)$$

where  $U_{\tilde{q}}$  is the mixing matrix from Exercise 22 and  $\omega_{\pm} = \frac{1}{2}(1 \pm \gamma_5)$  are the chiral projectors. Calculate the one-loop contribution to the quark self-energy  $i\Sigma^{\bar{q}q}(p)$  with momentum transfer  $p$ , induced by exchange of a SUSY particle which is mediated by the following diagram,



Parametrize the loop integrals by the following standard functions,

$$B_0(p^2, m_0, m_1) = \frac{(2\pi\mu)^{4-D}}{i\pi^2} \int d^D q \frac{1}{(q^2 - m_0^2 + i0) [(q+p)^2 - m_1^2 + i0]}, \quad (6)$$

$$B_{\mu}(p, m_0, m_1) = \frac{(2\pi\mu)^{4-D}}{i\pi^2} \int d^D q \frac{q_{\mu}}{(q^2 - m_0^2 + i0) [(q+p)^2 - m_1^2 + i0]} \quad (7)$$

$$= p_{\mu} B_1(p^2, m_0, m_1), \quad (8)$$

and decompose the self-energy into its vector-, axial vector-, and scalar parts according to the relation

$$\Sigma^{\bar{q}q}(p) = \not{p} \Sigma_V^{\bar{q}q}(p^2) + \not{p} \gamma_5 \Sigma_A^{\bar{q}q}(p^2) + m_q \mathbf{1} \Sigma_S^{\bar{q}q}(p^2). \quad (9)$$

Calculate the limit of large SUSY parameters (here equivalent to  $p^2, m_q \rightarrow 0$ ) analogously to Exercise 22.